

COMPARISON OF HADRONIC  
RESCATTERING CALCULATIONS  
OF ELLIPTIC FLOW AND HBT  
WITH MEASUREMENTS FROM RHIC

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OUTLINE

- 1) RESCATTERING CALCULATIONS
- 2) RESULTS FOR RADIAL AND ELLIPTIC FLOW
- 3) RESULTS FOR HBT

→ MOTIVATION ←

HOW FAR CAN ONE GET WITH A "PURE" (NAIVE)  
HADRONIC RESCATTERING SCENARIO ??

# RESCATTERING CALCULATION

USE HERMANN - BERTSCH METHOD FOR ISOSPIN-AVERAGED "PION-GAS" CALCULATION (PRC 51 (1995)), BUT INCLUDE BARYONS.

→  $\pi, p, \omega, \eta, \eta', K, K^*, \phi, N, \Delta, \Lambda$

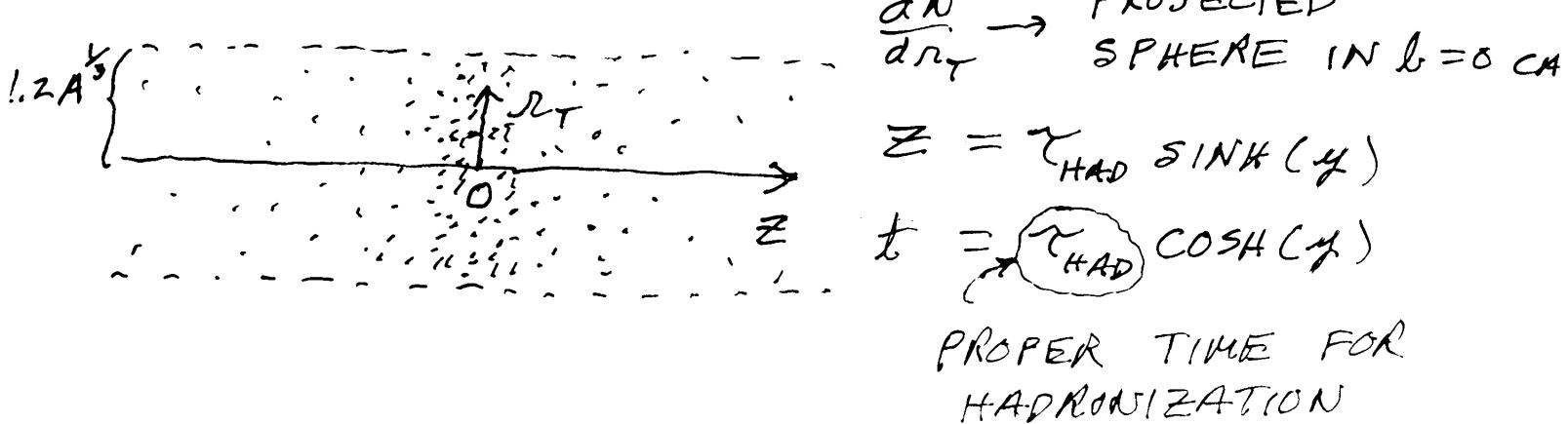
⇒ INITIAL STATE + RESCATTERING → FREEZEOUT

\* INITIAL STATE:  $T$  = TEMPERATURE PARAMETER

$$\frac{1}{m_T} \frac{d^2N}{dm_T dy} \propto \frac{m_T}{e^{m_T/T} \pm 1} f(y) \quad \text{RAPIDITY WIDTH}$$

WHERE  $f(y) = \begin{cases} e^{-y^2/2\sigma_i^2}, & i = \pi, K, p, \dots \\ \text{"FLAT" OR "PEAKED"}, & \text{BARYON} \end{cases}$  (MESONS)

→ GEOMETRY:



→ MULTIPICITIES:

CHOOSE  $\langle m_T \rangle$  TO EVOLVE TO EXPERIMENTAL  $\langle k_T \rangle$   
 (USE HELIOS p+p DATA TO HELP SET INITIAL  $\langle p_T \rangle$   
 RESONANCE FRACTIONS)

\* RESCATTERING:

"TUNE"  $T, g, \tau_{\text{had}}$  ⇌ ADJUSTABLE PARAMETERS

	(158 GeV Pb+Pb)	(130 GeV Au+Au)
<u>SPS</u>	<u>RHIC</u>	
T (MeV)	220	300
$\Sigma_{had}$ (fm/c)	1	1
$\sigma$	1.2	2.4

$v_2$  elliptic flow  
radial flow

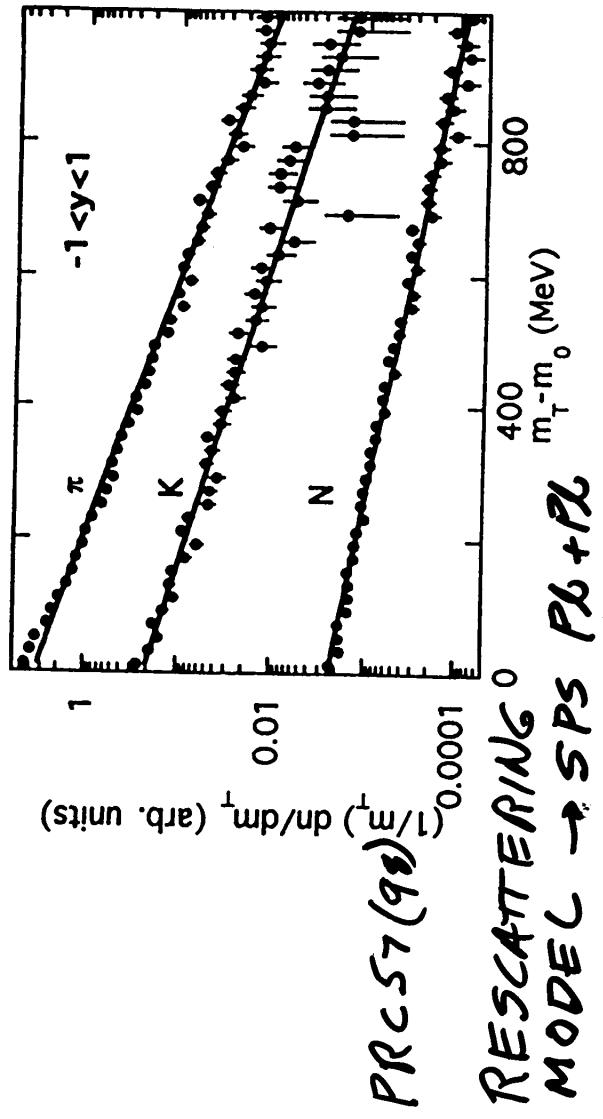
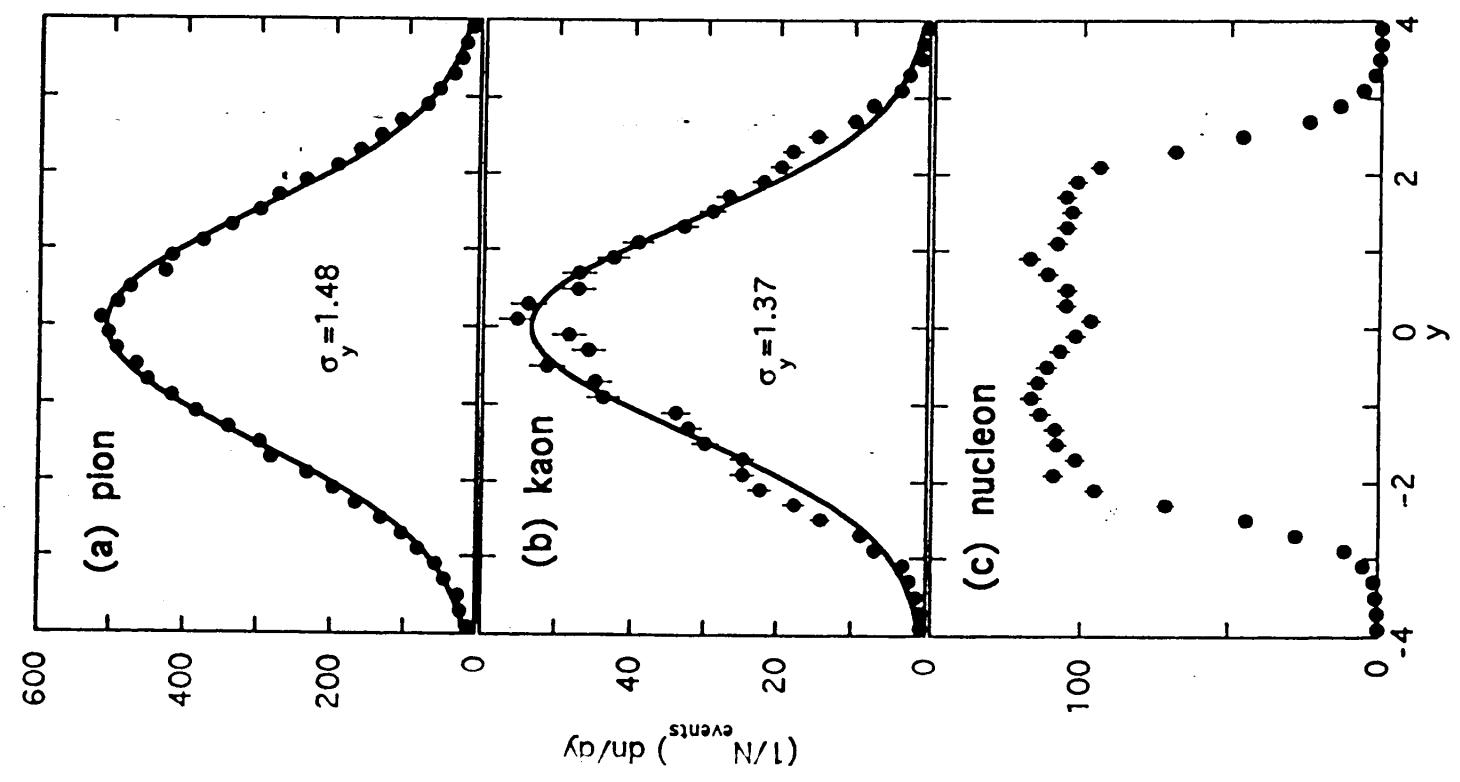


FIG. 6.  $m_T$  distributions at freeze-out from the rescattering calculation for pions, kaons, and nucleons at midrapidity. Normalizations are arbitrary. Fits to Eq. (1) are also shown.

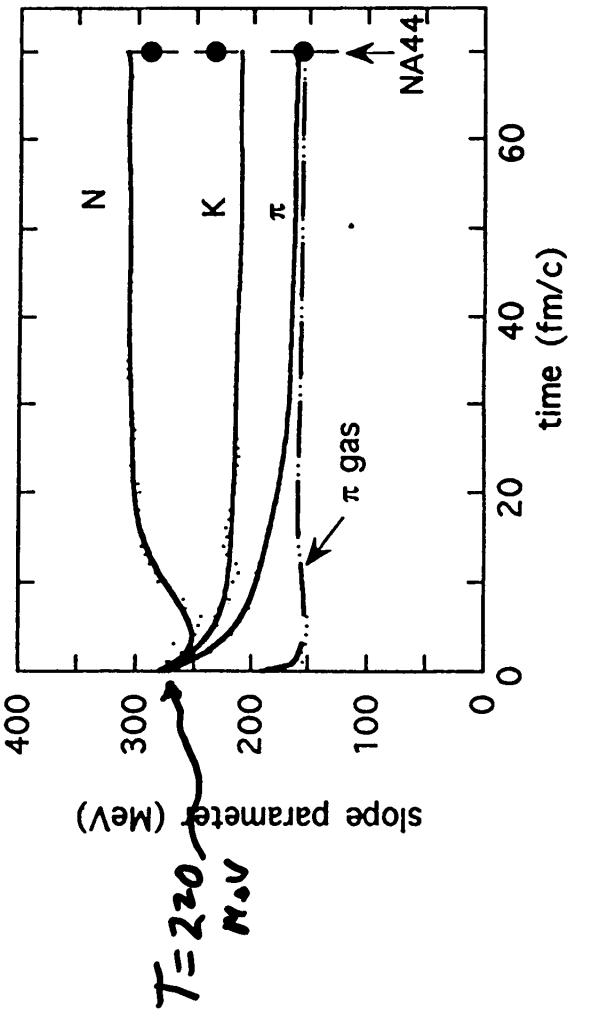


FIG. 5. Freeze-out rapidity distributions for (a) pions, (b) kaons, and (c) nucleons from the rescattering calculation. Fits to Eq. (2) for

FIG. 7. Time evolution of the slope parameters in the rescattering calculation for a full calculation (solid lines) and for a "pion gas" calculation (dashed line). Experimental freeze-out slope pa-

(30 GeV Au+Au RESCATT. CALC.

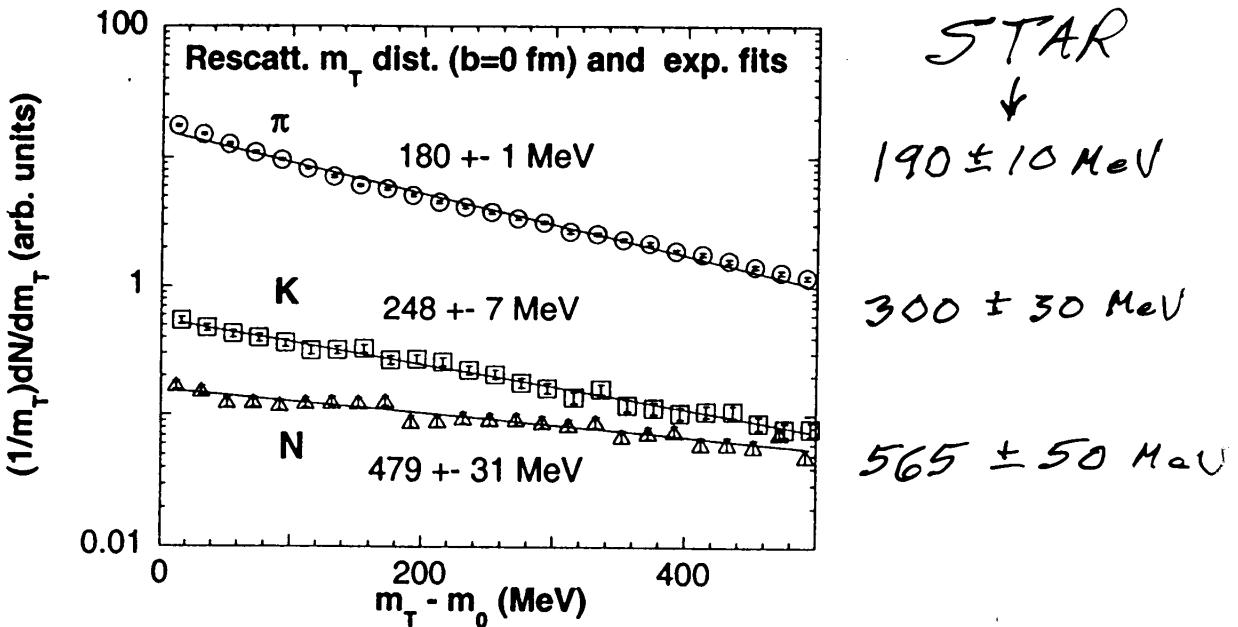


FIG. 2: Transverse mass distributions from the rescattering model. The lines are exponential fits to the distributions and the slope parameters are shown.

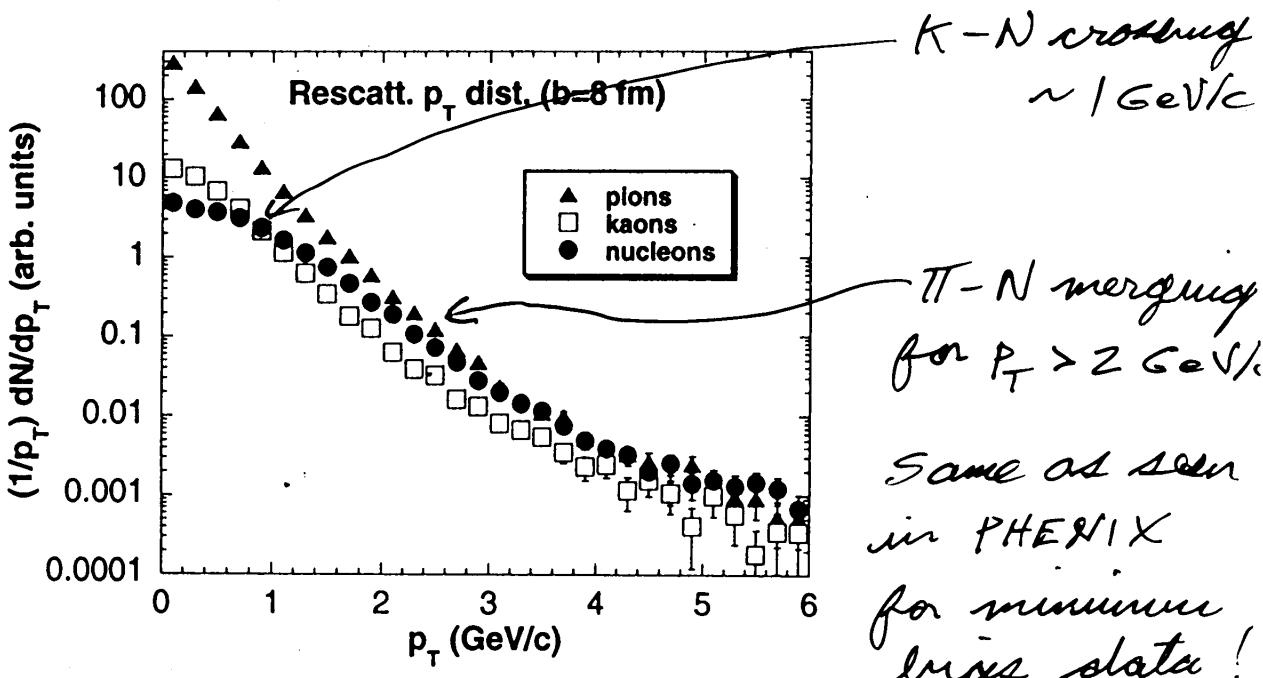


FIG. 3: Transverse momentum distributions for pions, kaons, and nucleons from a  $b = 8$  fm rescattering calculation.

Flattening out of  $\frac{dN}{d\eta}$  near  $\eta = 0$  due to transformation of Gaussian  $\frac{dN}{dy} \rightarrow \frac{dN}{d\eta}$

→ Flattening seen in PHOBOS

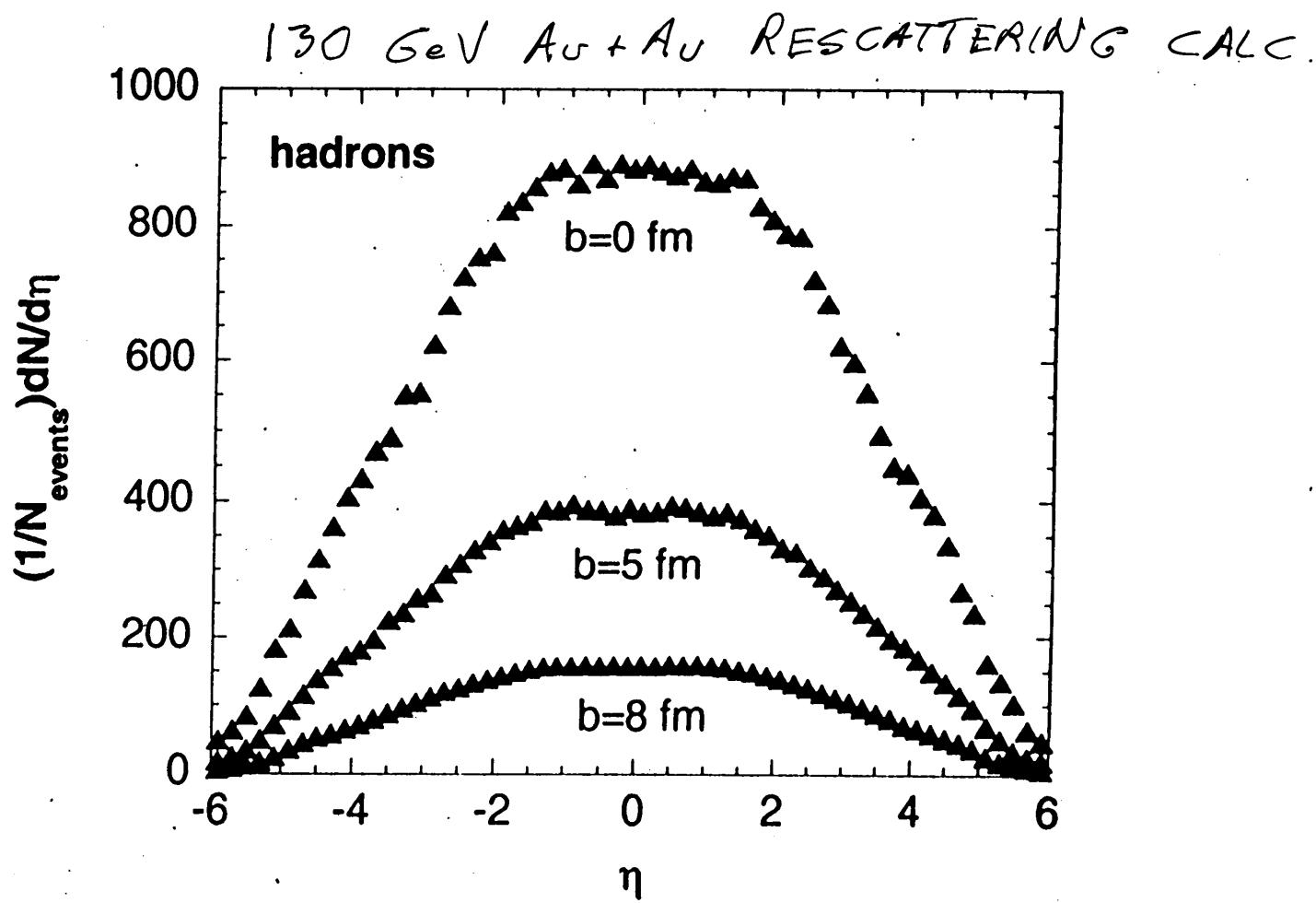
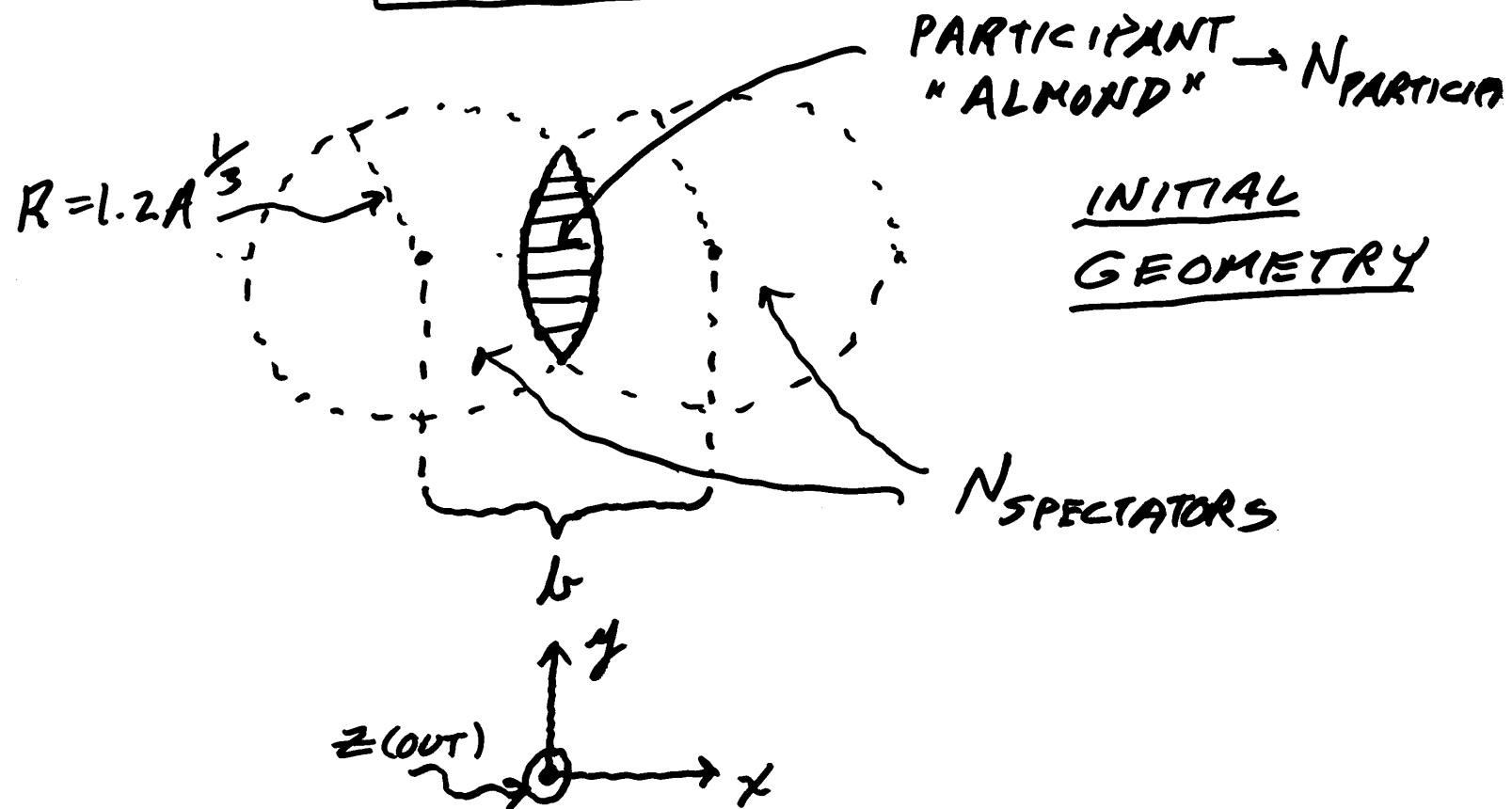


FIG. 1: Pseudorapidity distributions summed over pions, kaons, and nucleons from rescattering calculations for  $b = 0, 5$ , and  $8 \text{ fm}$ .

## $b > 0$ CALCULATIONS



FOR PARTICLE MULTIPLICITIES, SCALE  $\langle m_{b=0} \rangle_i$

$$\langle m_b \rangle_i = \frac{N_{\text{PARTICIPANTS}}}{N_{\text{PARTICIPANTS}} + N_{\text{SPECTATORS}}} \langle m_{b=0} \rangle_i$$

CALCULATE FOR  $b = 5\text{ fm}$ ,  $9\text{ fm}$

ELLIPTIC FLOW VARIABLE:

$$v_2 = \langle \cos 2\varphi \rangle$$

$$\text{WHERE } \varphi = \tan^{-1} \frac{P_y}{P_x}$$

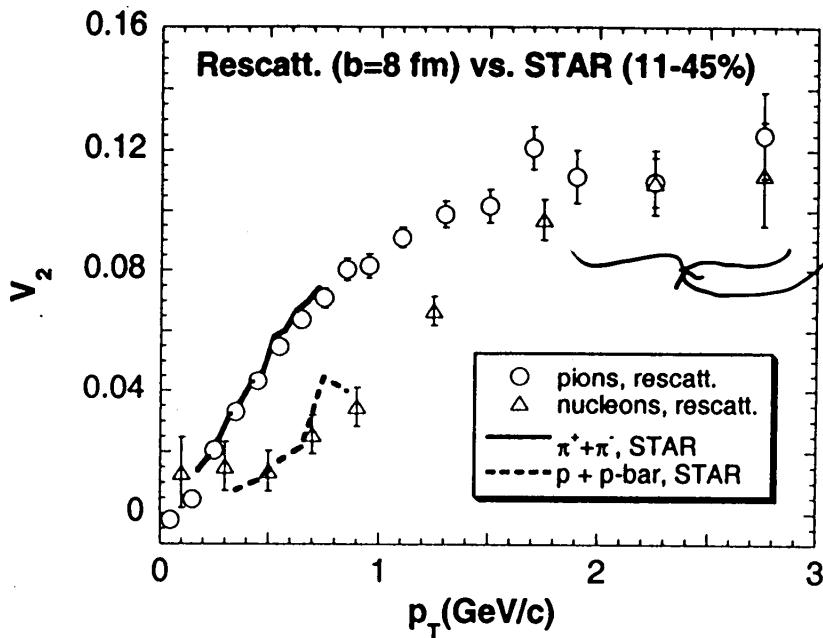


FIG. 4: Calculations of  $v_2$  from the rescattering model for  $b = 8$  fm for pions and nucleons compared with STAR measurements at 11-45% centrality. The plotted points with error bars are the rescattering calculations and the lines show the trends of the STAR measurements. Average errors on the STAR measurements are  $\leq 0.002$  for pions and 0.006 for protons+antiprotons.

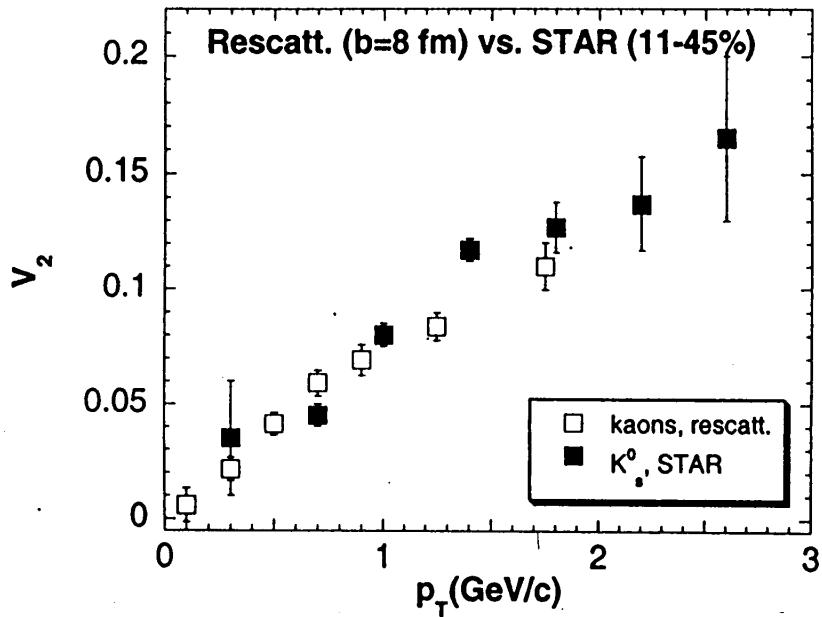


FIG. 5: Calculation of  $v_2$  from the rescattering model for kaons at  $b = 8$  fm compared with STAR measurements for  $K^0$ , at 11-45% centrality.

# HBT RESCATTERING CALCULATIONS

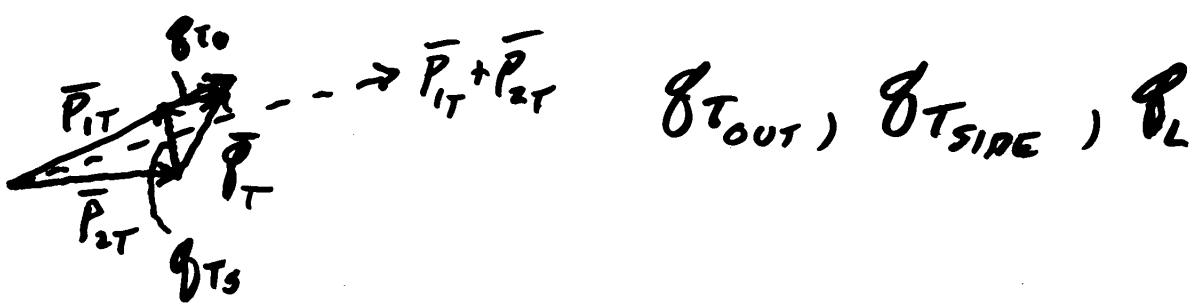
FOR HBT, IMPOSE SYMMETRIZATION ON  
PION PAIRS  $\{x_i, t_i, \bar{p}_i, m_i\}$   
 $\{x_j, t_j, \bar{p}_j, m_j\}$

$$W(i,j) = 1 + \cos[(\bar{P}_i - \bar{P}_j) \cdot (\bar{x}_i - \bar{x}_j) - (E_i - E_j)(t_i - t_j)]$$

BIN WEIGHTED PAIRS TO FORM CORRELATION FUNCT.

$$C(\bar{g}) = \frac{\sum_{i \neq j} \frac{dN}{d\bar{P}_i} \frac{dN}{d\bar{P}_j} W(i,j)}{\sum_{i \neq j} \frac{dN}{d\bar{P}_i} \frac{dN}{d\bar{P}_j}}, \quad \bar{g} = \bar{P}_i - \bar{P}_j$$

USE BERTSCH-PRATT VARIABLES (LCMS)



# FIT $C(\vec{g})$ WITH GAUSSIAN

$$C_G(\bar{g}) = A \left( 1 + \lambda e^{-g_{T_3}^2 R_{T_3}^2 - g_{T_0}^2 R_{T_0}^2 - g_c^2 R_c^2} \right)$$

$\Rightarrow \text{GET } \lambda, R_{T_3}, R_{T_0}, R_L$

# RESCATTERING AND SPS $\mu + \bar{\nu}$

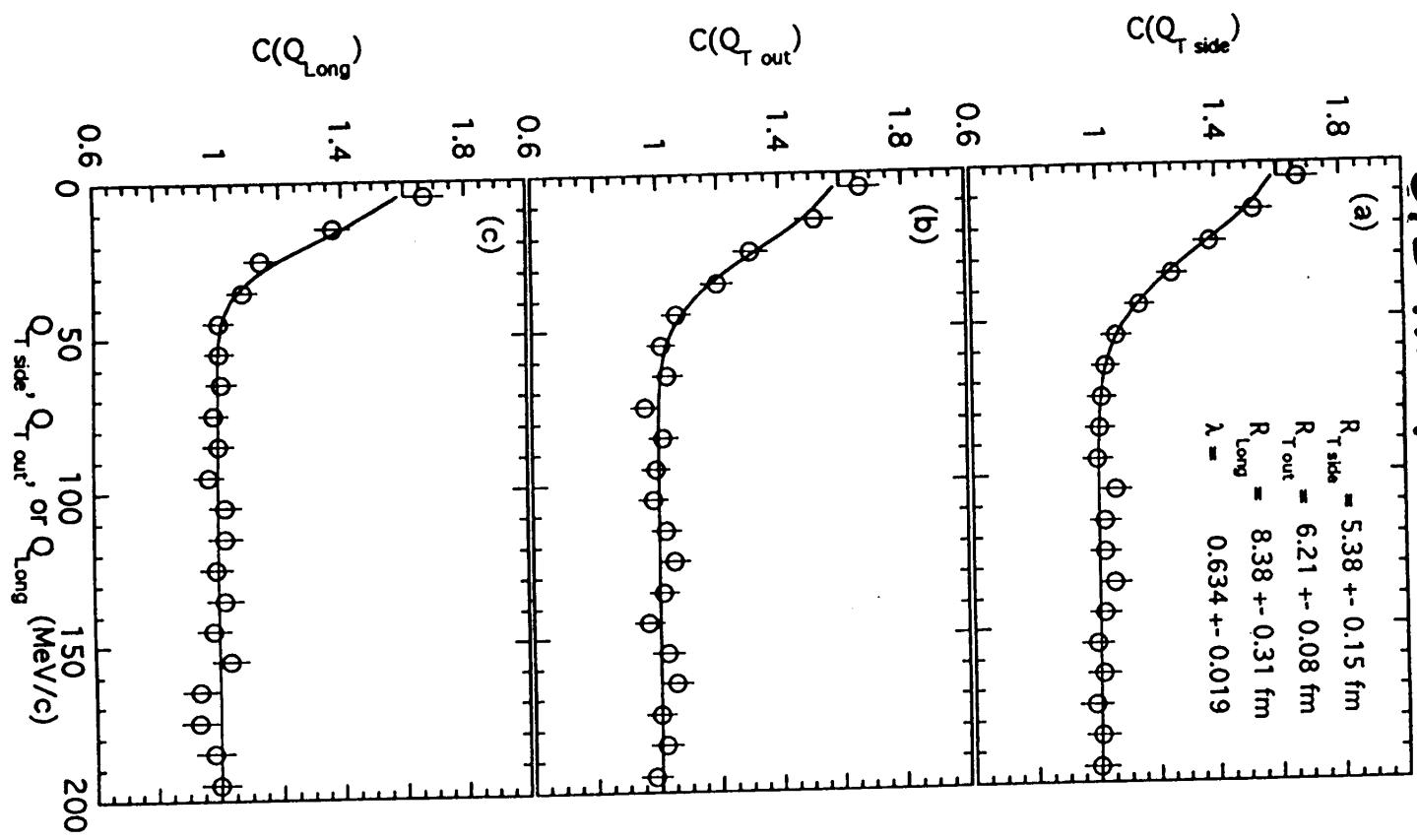


FIG. 8 Projections of the three-dimensional two-pion correla-

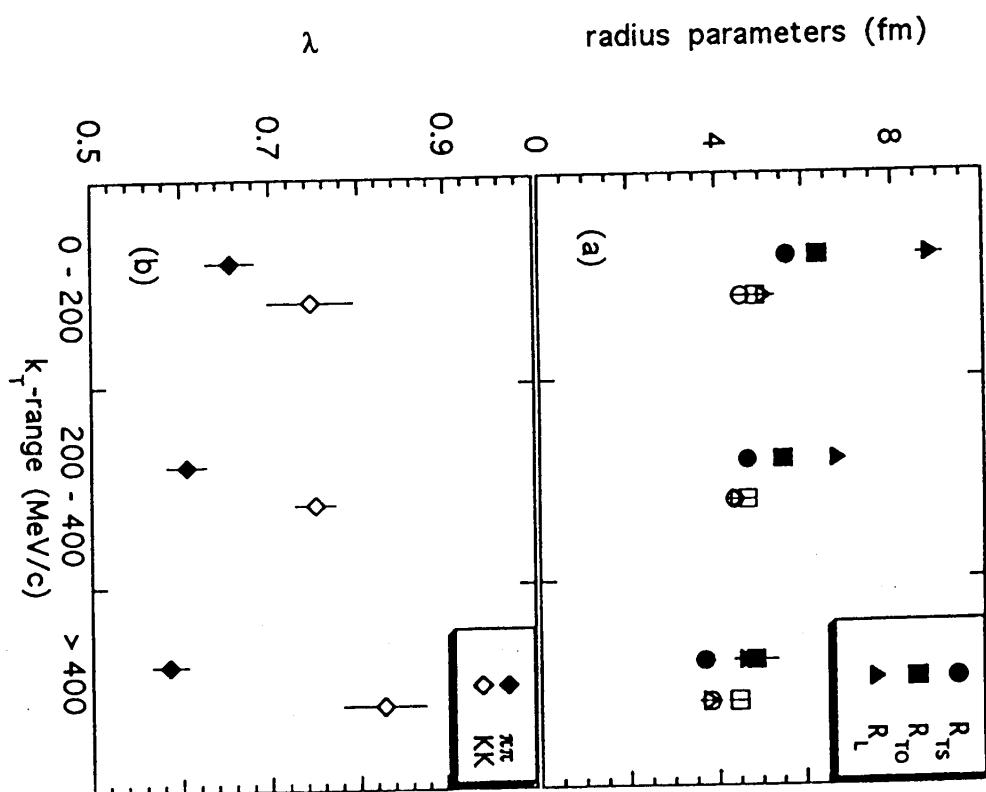


FIG. 9.  $k_T$  dependences of pion (solid symbols) and kaon (open symbols) (a) radius and (b)  $\lambda$  parameters extracted at midrapidity ( $-1 < y < 1$ ) from the present rescattering calculation.

fit with Eq. (9) has also been superimposed onto the Monte Carlo points, the extracted pion source parameters being displayed in the figure. The pion acceptance used for the calculation of this correlation function was  $-1 < y < 1$  and  $p_T < 400 \text{ MeV}/c$ , leading to an average  $k_T$  ( $k_T = |p_{T_1} + p_{T_2}|/2$ ) of about  $140 \text{ MeV}/c$ . The main differences between these source parameters and ones extracted from rescattering cal-

# RESCATTERING VS. NA49

→ 160 GeV/n Pb + Pb

57

CONSTRAINING A SIMPLE HADRO

HOMANIC, PRC 57 (1998)

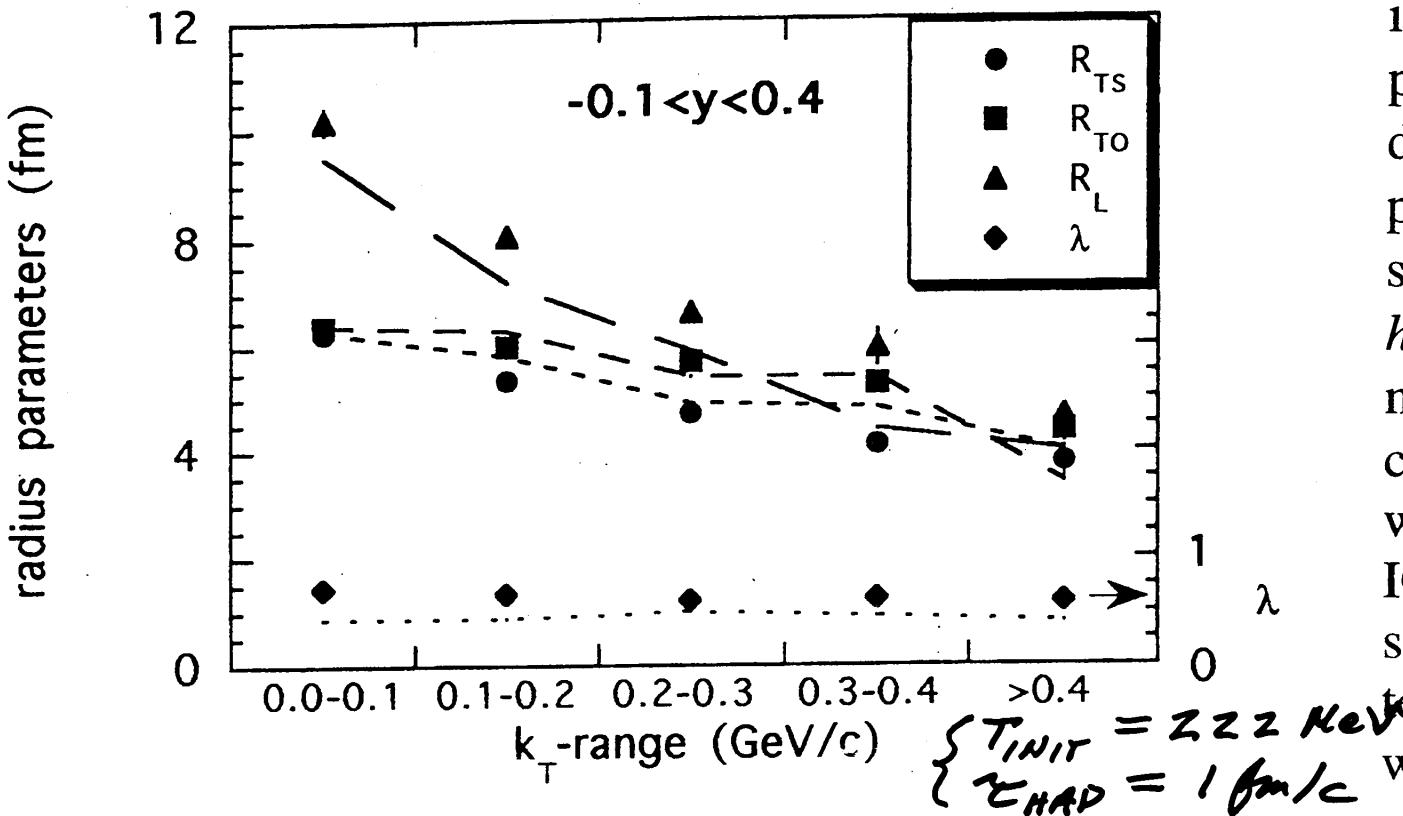


FIG. 11. Comparisons between calculations and NA49 measurements of the  $k_T$  dependence of the pion source parameters. The trends of the NA49 measurements are indicated by the dashed lines, which have similar identifications as in Fig. 10.

H. APPELSHÄUSER, PH.D. THESIS (1996)

dences of pion and kaon source parameters extracted at midrapidity ( $-1 < y < 1$ ) from the present rescattering calculation for central Pb+Pb collisions. Looking at the radius parameters first, there is a general trend for these parameters to decrease with increasing average  $k_T$  which is rather strong in the case of pions and weaker for kaons. Comparing the pion and kaon radius parameters at low  $k_T$ , the most striking difference seen between them is  $R_{long}$  is significantly larger

of about 140 MeV/c. The main difference between these source parameters and ones extracted from rescattering calculations without resonances [5] is seen in the  $\lambda$  parameter: without resonances  $\lambda \approx 1$  and with resonances  $\lambda < 1$ , in this case about 0.6. As has been pointed out previously [6], this reduction in  $\lambda$  is mainly due to the presence of long-lived resonances (e.g.,  $\eta$  and  $\eta'$ ). Figure 9 shows the  $k_T$  dependent

***PRC 57 (98) RESCATTERING MODEL  
SPS Pb + Pb***

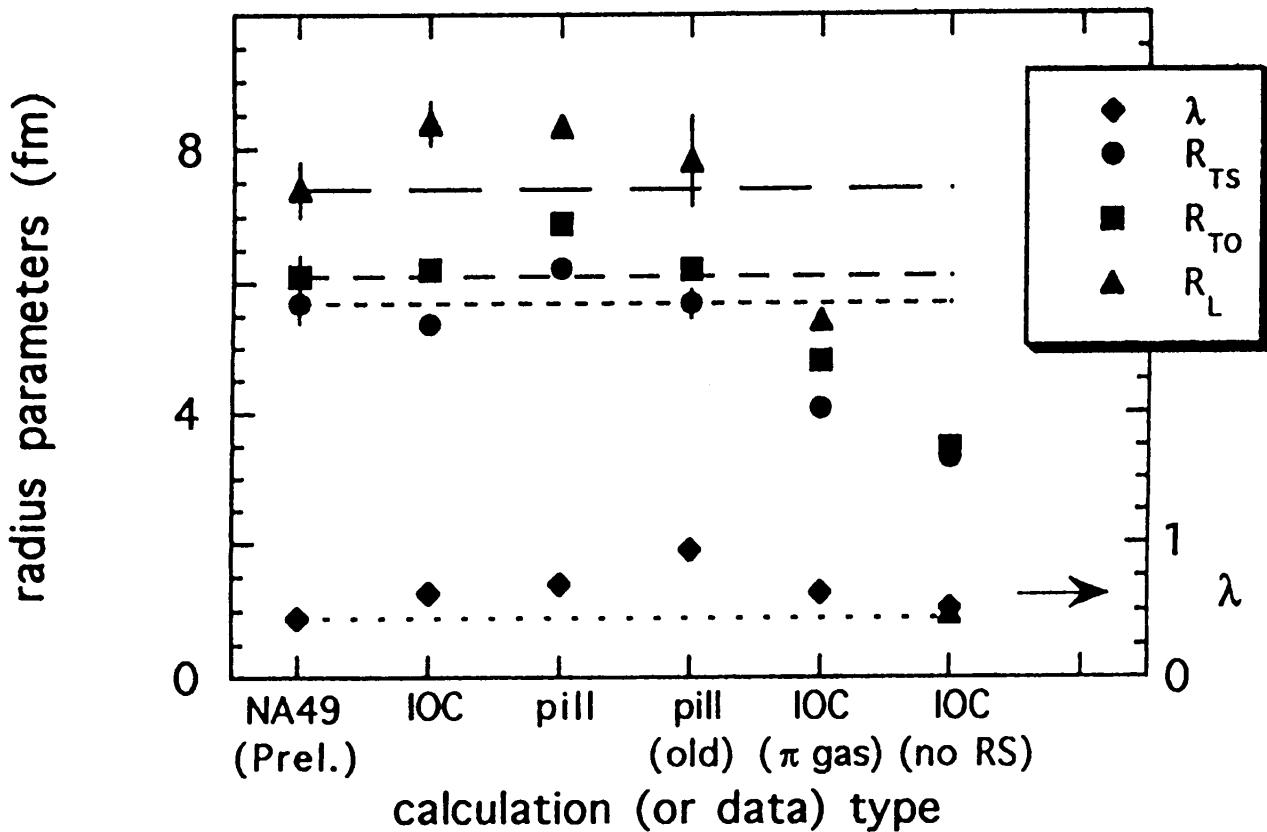


FIG. 10. Comparison between preliminary NA49 experimental pion interferometry results for Pb+Pb for the same average  $k_T$  and  $y$  range as used in Fig. 8 [10,11] and various rescattering model calculations. The dashed lines are projections of the experimental data points to allow an easier comparison with the calculations.

## COMPARISON OF RESCATTERING WITH STAR HBT

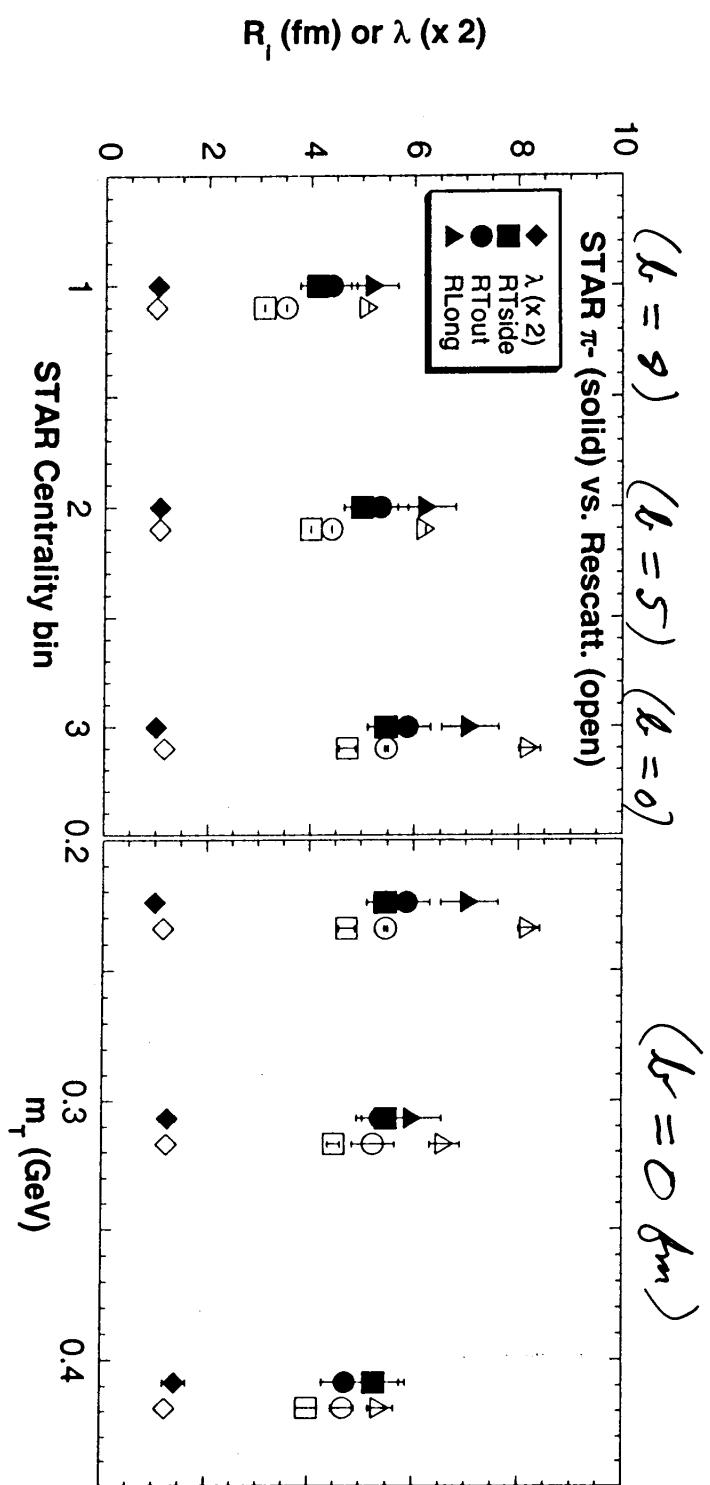
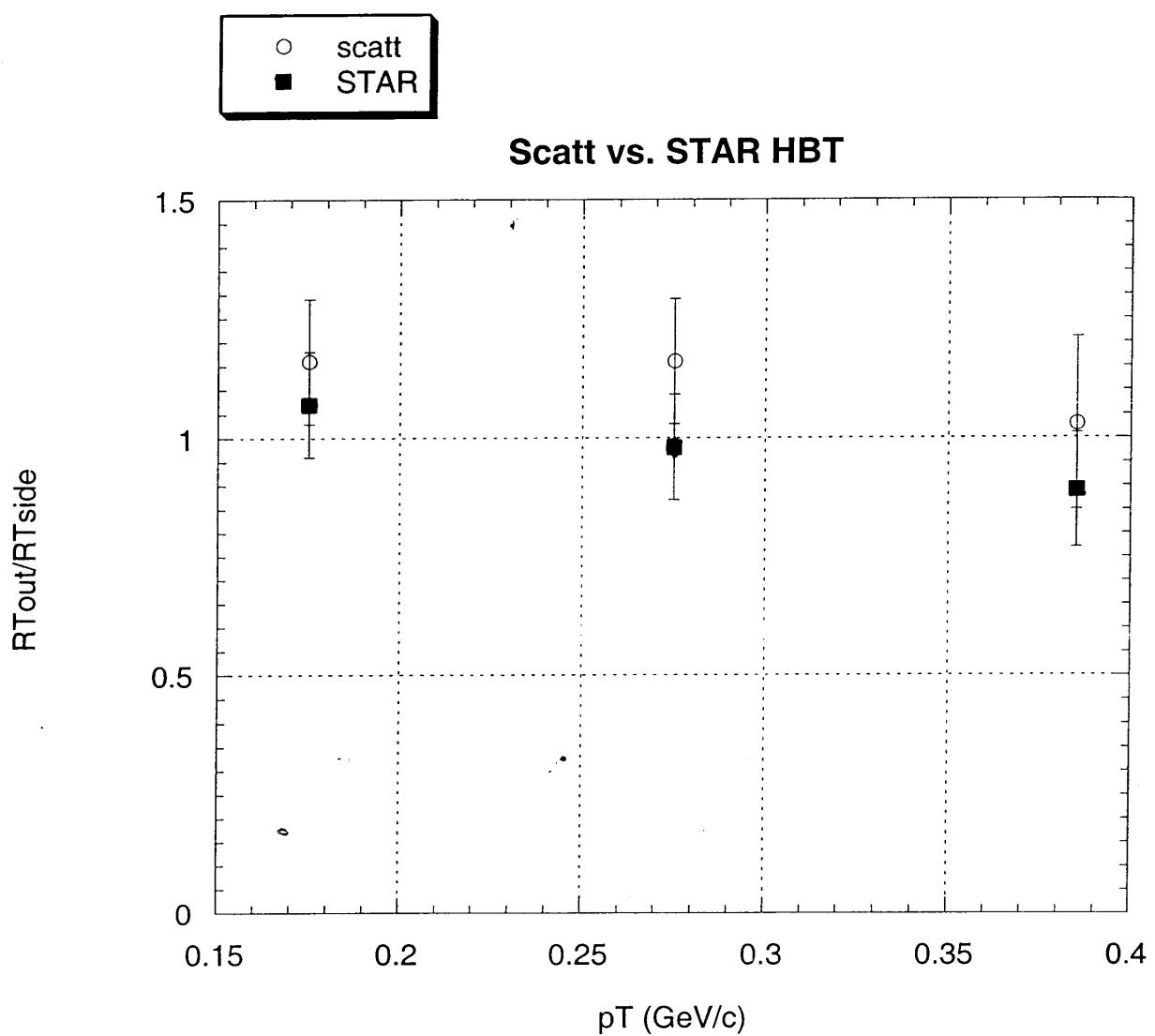
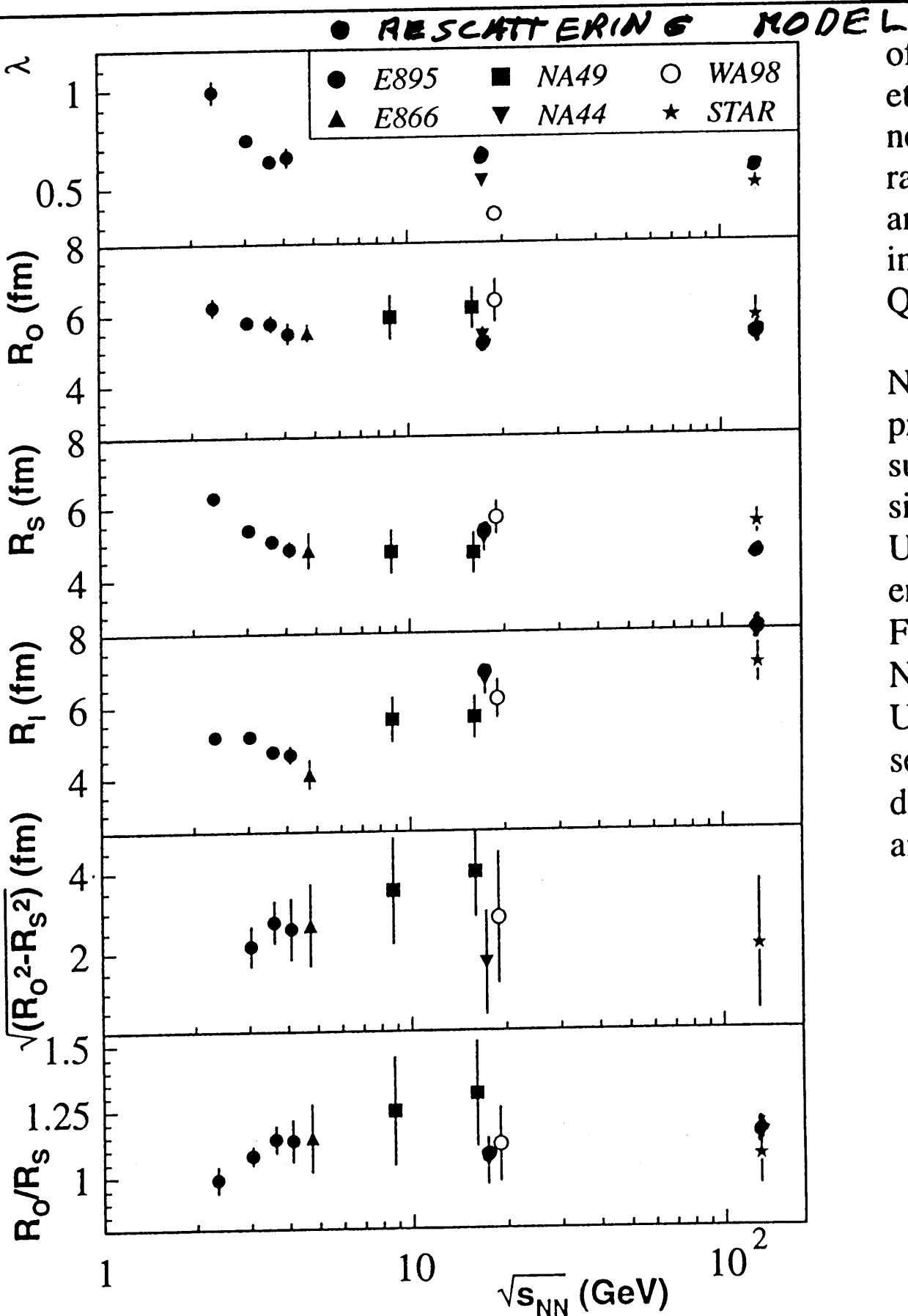


FIG. 3: Comparison of HBT source parameters from rescattering with STAR measurements as a function of centrality (see text) and  $m_T$ . The STAR measurements are the solid symbols and the rescattering calculations are the open symbols. The errors on the STAR measurements are statistical+systematic.

$b=0$  RESCATTERING vs. STAR HBT  
MODEL  
(CENTRAL TRIGGER)





13 The energy dependence of  $\pi^-$  HBT parameters for  
14  $\pi^- + \text{Au}$  and  $\pi^- + \text{Cu}$  collisions at  $v_{\text{rel}} = 0.15$  and  $p_T \approx$

of opaqueness in  
try excitation fu  
nearly two decade  
radii are observed  
are needed to com  
in emission time  
QGP formation.

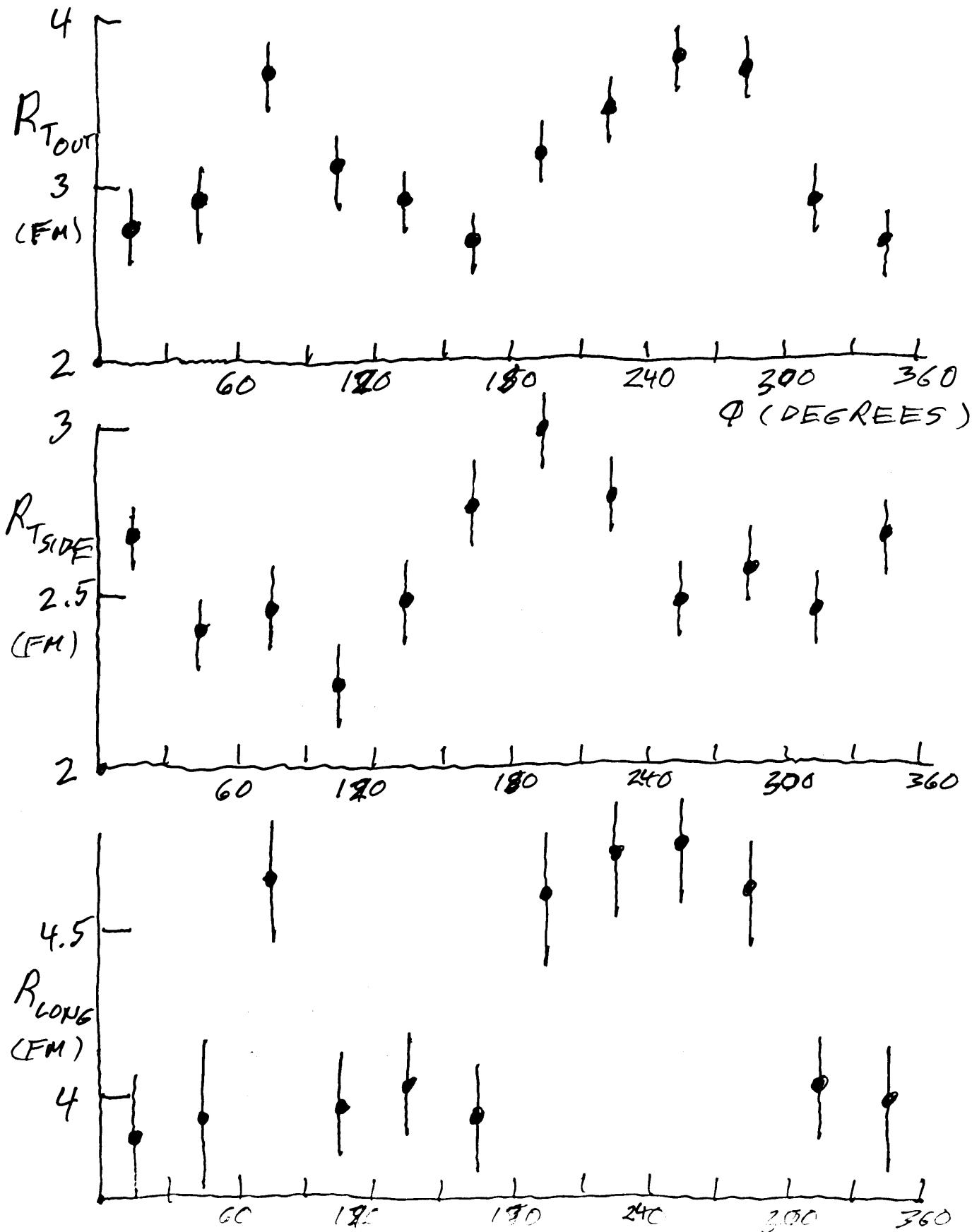
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and Technology.

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# HBT VS. $\phi$ FROM RESCATTERING $b = 8$ FM

(FROM 2 AM THIS MORNING  
CALCULATION)

(RHIC 130 GeV)



## CONCLUSIONS

- ⇒ THE RESCATTERING MODEL DOES A GOOD JOB AT DESCRIBING
- \* RADIAL FLOW
  - \* ELLIPTIC FLOW
  - \* HBT
- FOR RHIC  $\sqrt{s} = 130 \text{ GeV}$  Au+Au DATA
- ⇒ SUGGESTS THE PRE-HADRON STATE OF THE COLLISION IS SHORT-LIVED ( $\tau \sim 1 \text{ fm/c}$ ) AND "HOT" ( $T \sim 300 \text{ MeV}$ )